

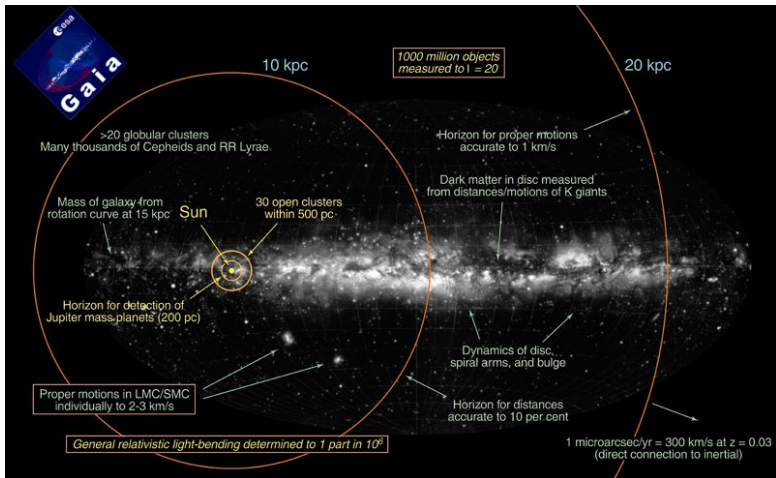
Time Transfer functions as tool to validate light propagation solutions for space astrometry

S. Bertone*, O. Minazzoli, M. Crosta, C. Le Poncin-Lafitte,
A. Vecchiato, M.-C. Angonin

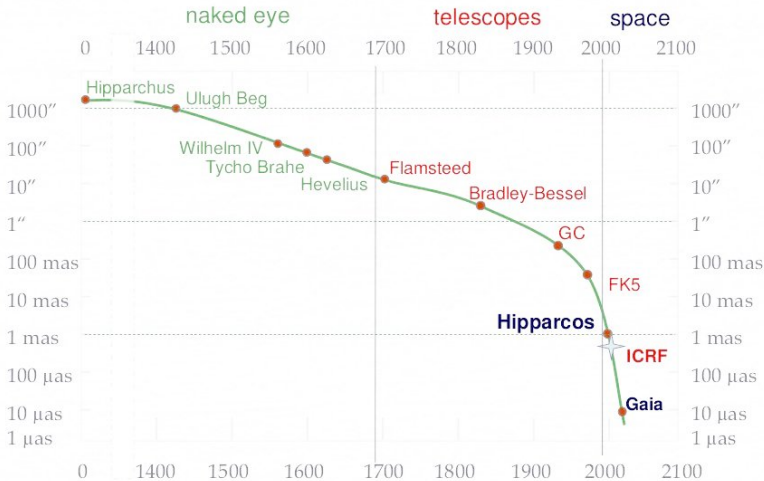
* SYRTE/Obspm - INAF/Università di Torino

June 07, 2013
Journées de la SF2A

- 1 Astrometry in Gaia's age
- 2 The astrometric core solution
- 3 Relativistic models for light propagation
- 4 A cross-check procedure
- 5 Conclusions



Schematic diagram showing the distances out to which Gaia will contribute to our knowledge of the Galaxy.
Image: ESA



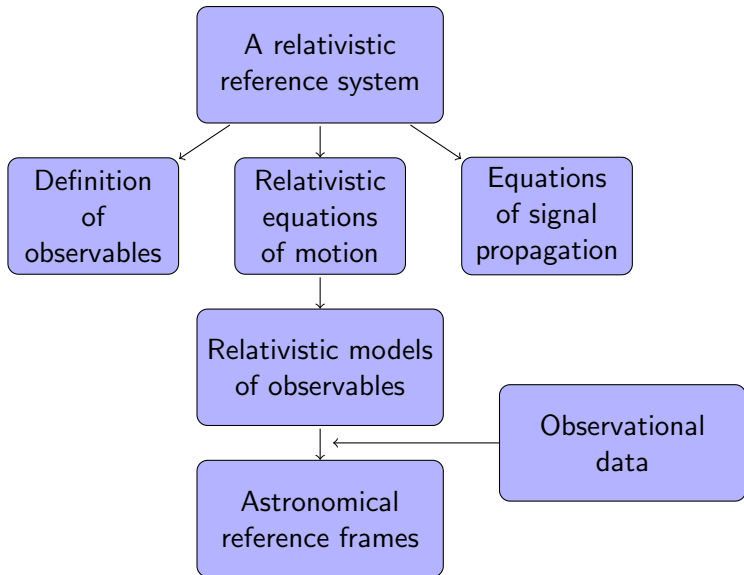
Accuracy of astrometric observations VS year. Image : S. Klioner, Porto 2011

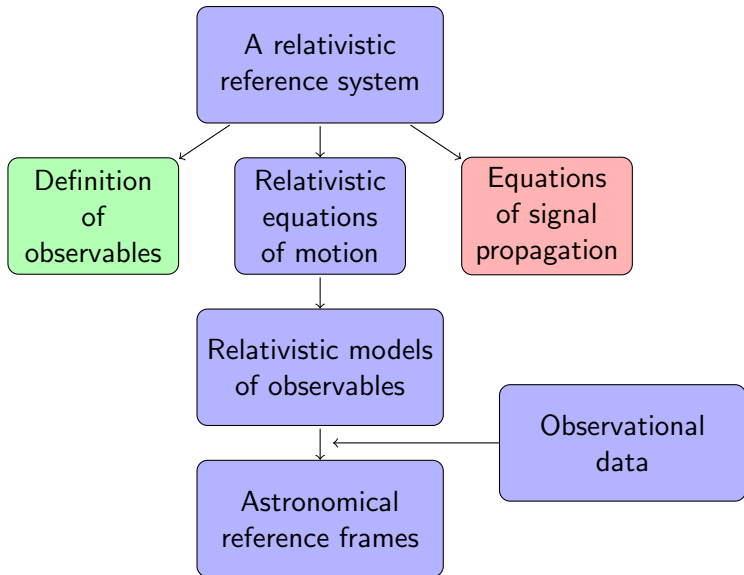
Relativistic deflection by Solar
System planets : $10 - 10^4 \mu\text{as}$

Relativistic light deflections in the Solar System

Body	Monopole		Quadrupole	
	grazing mas	χ $\delta\theta = 1 \mu\text{as}$	grazing μas	χ $\delta\theta = 1 \mu\text{as}$
Sun	17,000	180°		
Mercury	0.083	0.15°		
Venus	0.49	4.5°		
Mars	0.12	0.4°		
Jupiter	16.3	90°	240	$8 R_J$
Saturn	5.8	17°	95	$4 R_S$
Uranus	2.1	1.2°	8	$2 R_U$
Neptune	2.5	0.9°	10	$2 R_N$

Mignard, Klioner , Gaia : Relativistic modelling and testing, 2009





Astrometric Global Iterative Solution (AGIS)

Global Sphere Reconstruction (GSR)

Procedure

Definition
of the
observable

Light
propagation

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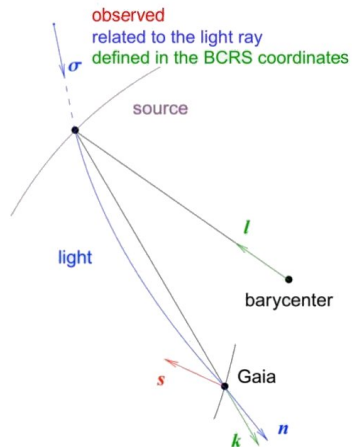
RAMOD

General RELativistic Model (GREM)

Based on IAU reference systems

$$s \leftrightarrow n \leftrightarrow \sigma \leftrightarrow k \leftrightarrow l, \pi$$

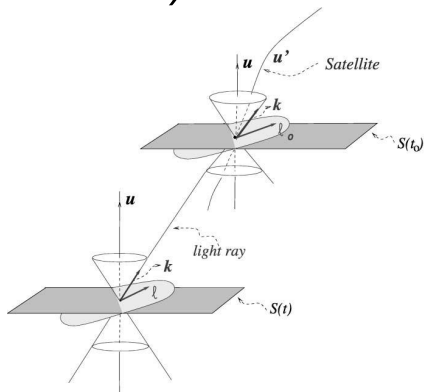
- (1) aberration
- (2) gravitational deflection
- (3) coupling to finite distance
- (4) parallax



General structure of the General RELativistic Model (GREM). Image: S. Klioner, PRD 2003

Relativistic Astrometric MODel (RAMOD)

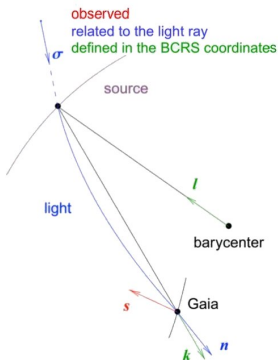
- Based on a measurement protocol
- local barycentric observer u
- $\bar{\ell} =$ local line of sight of the fiducial observer u



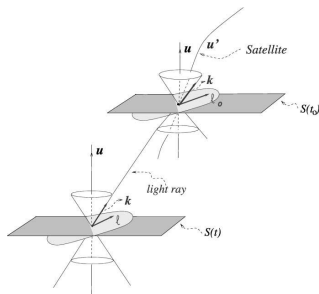
General structure of the Relativistic Astrometric Model (RAMOD). Image: M. Crosta, Porto 2011

$$\frac{d\bar{\ell}^k}{d\zeta} + \bar{\ell}^i \bar{\ell}^j \left(\partial_i h_{kj} - \frac{1}{2} \partial_k h_{ij} \right) + \frac{1}{2} \bar{\ell}^k \bar{\ell}^i \partial_i h_{00} - \frac{1}{2} \partial_k h_{00} + \mathcal{O}(h^2) = 0$$

How to relate their results?



General structure of the General Relativistic Model (GREM). Image: S. Klioner, PRD 2003



General structure of the Relativistic Astrometric Model (RAMOD). Image: M. Crosta, Porto 2011

Previous studies : aberration (M. Crosta and A. Vecchiato 2010), geodesic equations (M. Crosta 2011)

Time Transfer Functions (TTF)

$$\mathcal{T}_r(\mathbf{x}_A, t_B, \mathbf{x}_B) = \frac{R_{AB}}{c} + \sum_{n=1}^{\infty} G^n \Delta_r^{(n)}(\mathbf{x}_A, t_B, \mathbf{x}_B, g_{\mu\nu})$$
$$\hat{k}_i^{x_B} = \frac{k_i^{x_B}}{k_0^{x_B}} = -c \frac{\partial \mathcal{T}_r}{\partial x_B^i} \left[1 - \frac{\partial \mathcal{T}_r}{\partial t_B} \right]^{-1}$$

Le Poncin-Lafitte et al. 2004, Teyssandier & Le Poncin-Lafitte 2008; T. 2012.

- Δ_r at any order in general static, spherically symmetric space-times
- without integrating the whole set of geodesic equations
- well adapted to a ray emitted and observed at points both at a finite distance x_A et x_B
- definition of the astrometric observable within the formalism (Bertone and Le Poncin-Lafitte 2012)

Time Transfer Functions (TTF)

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Le Poncin-Lafitte et al. 2004, Teyssandier & Le Poncin-Lafitte 2008; T. 2012.

We propose to :

extract \mathcal{T} and \hat{k}_i from GREM and RAMOD

TTF formalism in closed form

(S. Bertone and C. Le Poncin Lafitte, Memorie SAI 2012)

$$\mathcal{T}_r(\mathbf{x}_A, t_B, \mathbf{x}_B) = \frac{R_{AB}}{c} + \frac{1}{c} \Delta_r(\mathbf{x}_A, t_B, \mathbf{x}_B) + \mathcal{O}(c^{-5})$$

$$\left(\widehat{k}_i\right)_B \approx N_{AB}^i + \frac{\partial \Delta_r}{\partial x_B^i} + N_{AB}^i \frac{\partial \Delta_r}{\partial x_B^0}$$

TTF formalism in closed form

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$$\Delta_r = \frac{1}{2} R_{AB} \int_0^1 \left[h_{00}^{(2)} + \frac{2}{c} N_{AB}^i h_{0i}^{(3)} + N_{AB}^i N_{AB}^j h_{ij}^{(2)} \right]_{z_-^\alpha(\lambda)} d\lambda$$

TTF for a time-dependent metric

(S. Bertone et al., 2013)

$$\text{PPN metric : } h_{00} = \frac{2G}{c^2} \sum_P \frac{\mathcal{M}_P}{R_P(t, \mathbf{x})}, \quad h_{0i} = -(1 + \gamma)h_{00}\beta_P^i(t), \quad h_{ij} = \delta_{ij}\gamma h_{00}$$

$$\mathbf{R}_P = \mathbf{R}_{PB}^* - \lambda R_{AB}(\mathbf{N}_{AB} - \beta_P)$$

$$\begin{aligned} \mathcal{T}_r(\mathbf{x}_A, t_B, \mathbf{x}_B) &= \frac{R_{AB}}{c} + (\gamma + 1) \frac{G}{c^2} \sum_P \mathcal{M}_P \left[1 - \beta_P(t_C) \cdot \mathbf{N}_{AB} \right] \\ &\quad \times \ln \left[\frac{R_{PA} - \mathbf{R}_{PA} \cdot \mathbf{N}_{AB} - \beta_P(t_C) \cdot (\mathbf{R}_{PA} - \mathbf{N}_{AB} R_{PA})}{R_{PB} - \mathbf{R}_{PB} \cdot \mathbf{N}_{AB} - \beta_P(t_C) \cdot (\mathbf{R}_{PB} - \mathbf{N}_{AB} R_{PB})} \right] \end{aligned}$$

$$\begin{aligned} (\hat{\mathbf{k}}_i)_B &= -N_{AB}^i + (\gamma + 1) \frac{G}{c^2} \sum_P \frac{\mathcal{M}_P}{R_{AB} R_{PB}} \left[R_{PB}^2 g_P^2 - (\mathbf{R}_{PB} \cdot \mathbf{g}_P)^2 \right] \\ &\quad \times \left\{ g_P N_{AB}^i \left[(\mathbf{R}_{PB} \cdot \mathbf{N}_{AB}) (\mathbf{R}_{PB}^2 - R_{PA} R_{PB} - R_{AB} \mathbf{R}_{PB} \cdot \beta_P(t_C)) \right. \right. \\ &\quad \left. \left. - R_{PB}^2 R_{AB} g^2 \right] + R_{PB}^i g_P^2 [R_{PB} R_{PA} - R_{PB}^2 + R_{AB} \mathbf{R}_{PB} \cdot \mathbf{g}_P] \right. \\ &\quad \left. + \beta_P^i(t_C) R_{PB} [(R_{PA} - R_{PB})(\mathbf{R}_{PB} \cdot \mathbf{N}_{AB}) + R_{PB} R_{AB}] \right\} \\ &\quad + (\gamma + 1) \frac{G}{c^2} \sum_P \mathcal{M}_P \frac{\beta_P^i(t_C) - N_{AB}^i \beta_P(t_C) \cdot \mathbf{N}_{AB}}{R_{AB} g_P} \ln \frac{g_P R_{PB} + \mathbf{R}_{PB} \cdot \mathbf{g}_P}{g_P R_{PA} + \mathbf{R}_{PA} \cdot \mathbf{g}_P} \\ &\quad + \mathcal{O}(c^{-4}). \end{aligned}$$

TTF for a static metric

$$\text{Static metric : } h_{00} = \frac{2G}{c^2} \sum_P \frac{\mathcal{M}_P}{R_P(t, \mathbf{x})}, \quad h_{0i} = 0, \quad h_{ij} = \delta_{ij} \gamma h_{00}$$

$$\mathbf{R}_P = \mathbf{x}_\gamma - \mathbf{x}_P = \mathbf{R}_{PB} - \lambda R_{AB} \mathbf{N}_{AB}$$

$$\mathcal{T}_r(\mathbf{x}_A, t_B, \mathbf{x}_B) = \frac{R_{AB}}{c} + (\gamma + 1) \frac{G}{c^2} \sum_P \mathcal{M}_P \ln \left[\frac{R_{PA} - \mathbf{R}_{PA} \cdot \mathbf{N}_{AB}}{R_{PB} - \mathbf{R}_{PB} \cdot \mathbf{N}_{AB}} \right]$$

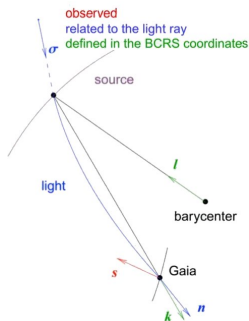
$$\begin{aligned} (\widehat{k}_i)_B &= -N_{AB}^i + (\gamma + 1) \frac{G}{c^2} \sum_P \frac{\mathcal{M}_P}{R_{AB} R_{PB} [R_{PB}^2 g_P^2 - (\mathbf{R}_{PB} \cdot \mathbf{N}_{AB})^2]} \\ &\times \left\{ N_{AB}^i \left[(\mathbf{R}_{PB} \cdot \mathbf{N}_{AB}) (R_{PB}^2 - R_{PA} R_{PB}) \right. \right. \\ &\quad \left. \left. - R_{PB}^2 R_{AB} \right] + R_{PB}^i \left[R_{PB} R_{PA} - R_{PB}^2 + R_{AB} \mathbf{R}_{PB} \cdot \mathbf{N}_{AB} \right] \right\} \\ &+ \mathcal{O}(c^{-4}). \end{aligned}$$

GREM < - > TTF

$$\mathcal{T} = t_B - t_A = \Delta t_{AB}$$

$$\hat{k}_i = \frac{k_i}{k_0} = \frac{g_{ij}k^j + g_{0i}k^0}{g_{00}k^0 + g_{0i}k^i}$$

$$\approx -\frac{\dot{x}^i}{c} - 2h_{00}\sigma^i - (\delta_{ij} + \sigma^i\sigma^j)h_{0j}$$



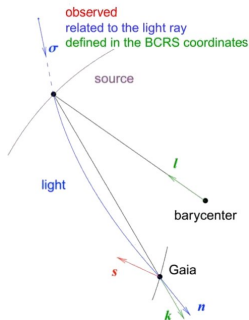
S. Klioner and S. Kopeikin 1992,
S. Klioner 2003

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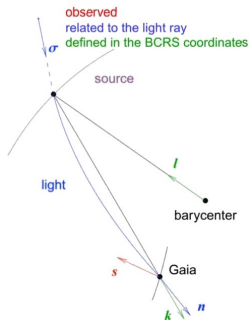
$$x^i(t) = x^i(t_B) + c\sigma^i(t - t_B) + \Delta x^i(t, x^i, t_B, x_B^i)$$

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$$x^i(t) = x^i(t_B) + c\sigma^i(t - t_B) + \Delta x^i(t, x^i, t_B, x_B^i)$$

$$\frac{\dot{x}^i(t, \mathbf{x})}{c} = \sigma^i + \frac{\Delta \dot{x}^i(t, \mathbf{x})}{c}$$

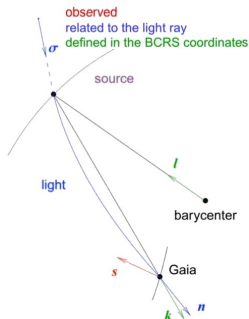
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$$\mathbf{x}(\mathbf{x}_B, \boldsymbol{\sigma}, \Delta t) = \mathbf{x}_A$$



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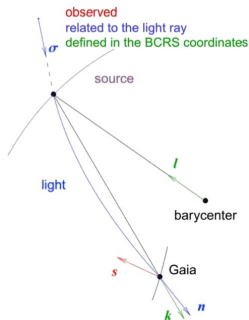
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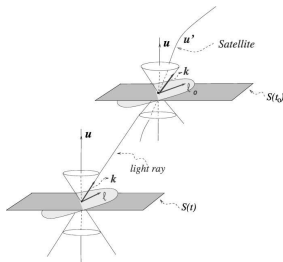
$$\mathbf{x}^i(t_A) = \mathbf{x}^i(t_B) + c\sigma^i\Delta t_{AB} + \Delta x^i(\Delta t_{AB}, R_{AB}^i)$$

$$\frac{\dot{x}^i(t, \mathbf{x})}{c} = \sigma^i + \frac{\Delta \dot{x}^i(t, \mathbf{x})}{c}$$

RAMOD < - > TTF

$$\mathcal{T} = t_B - t_A = \Delta t_{AB}$$

$$\bar{\ell}^i = -\frac{k^i}{u^0 k_0}$$

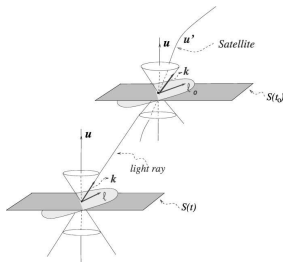


F. de Felice et al. 2004,
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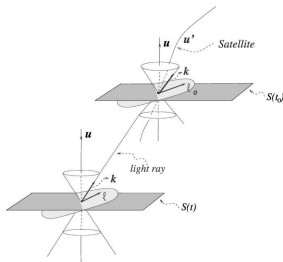
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$$u^0 \equiv \frac{cdt}{d\zeta} = \frac{1}{\sqrt{-g_{00}}}$$

RAMOD < - > TTF

$$\mathcal{T} = t_B - t_A = \Delta t_{AB}$$

$$\hat{k}_i^B = -\bar{\ell}_B^i \sqrt{g_{00}(x_B)} g_{ij}(x_B)$$



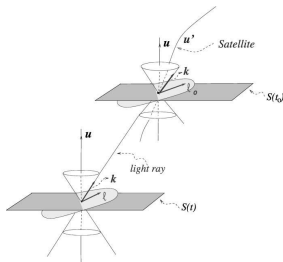
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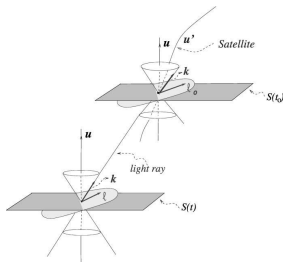
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$$\bar{\ell}_B^i = \frac{x_B^i - x_A^i}{\Delta\zeta_{AB}} + \frac{G}{c^2} \Delta\bar{\ell}(\mathbf{x}_A, \mathbf{x}_B, \Delta\zeta_{AB}, \mathbf{x}_P) + O(Gc^{-3})$$

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$$\mathcal{T} = t_B - t_A = \Delta t_{AB}$$

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F. de Felice et al. 2004,
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$$u^0 \equiv \frac{cdt}{d\zeta} = \frac{1}{\sqrt{-g_{00}}}$$

$$\Rightarrow c\Delta t_{AB} \approx \Delta\zeta_{AB} + \frac{G}{c^2} \Delta t_{AB}^{(2)}(\mathbf{x}_A, \mathbf{x}_B, \Delta\zeta_{AB}, \mathbf{x}_P)$$

$$\bar{\ell}_B^i = \frac{x_B^i - x_A^i}{\Delta\zeta_{AB}} + \frac{G}{c^2} \Delta\bar{\ell}(\mathbf{x}_A, \mathbf{x}_B, \Delta\zeta_{AB}, \mathbf{x}_P) + O(Gc^{-3})$$

Conclusions

- Absolute high-precision astrometry needs independent verifications
- More than one model of relativistic light propagation to interpret experimental data
- We propose a procedure to relate and understand them using the TTF formalism

Perspectives

- Development of a GSR-TTF

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